Two-Colour Nematicons in the Local Response Regime

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Nematic Liquid Crystals



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Experimental Results



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Governing Equations



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Variational Equations



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NLCs and Nematicons



Nematic liquid crystal: rod-like molecules.

Molecules tend to re-orientate in direction of electric field.

Coherent, polarised optical beam.

Beam intensity reduces angle of nematic molecules w.r.t. electric field.

Beam self focuses.

- **Different** wavelengths \Rightarrow different degrees of birefringence and dispersion.
- Cross-phase modulation (CPM) negates this effect.
- Self-localised VS forms.



Governing Equations

Single Nematicon Governing Equations

NLS-like equation for electric field of the light.

Forced heat equation for the director.

Two-Colour Nematicon Governing Equations

CNLS-like equation for each colour

$$iu_t + \frac{D_1}{2}\nabla^2 u + Au\sin 2\theta = 0$$
$$iv_t + \frac{D_2}{2}\nabla^2 v + Bv\sin 2\theta = 0.$$

Director equation

$$q\sin 2\theta - 2A|u|^2\cos 2\theta - 2B|v|^2\cos 2\theta = \nu\nabla^2\theta$$

ν large and ν small

 ν : ratio of elastic energy of nematicon to energy of electric field. Nonlocality parameter.



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- $\nu \rightarrow \infty$ for nonlocal crystals.
 - Low power input beam. Studied experimentally extensively.
 - Nonlocality stops collapse.
- $\nu \rightarrow 0$ for local crystals.
 - **S**trongly illuminated or heated NLCs.
 - Stability caused by saturation of NLC director motion.



Consider director equation for ν small <-

Rearrange

$$\tan 2\theta = \frac{2\left(A|u|^2 + B|v|^2\right)}{q}$$

Assuming radial symmetry, the governing equations are

$$iu_t + \frac{D_1}{2}u_{rr} + \frac{D_1}{2r}u_r + Au\left(\frac{2(A|u|^2 + B|v|^2)}{\sqrt{q^2 + 4(A|u|^2 + B|v|^2)}}\right) = 0$$
$$iv_t + \frac{D_2}{2}v_{rr} + \frac{D_2}{2r}v_r + Bv\left(\frac{2(A|u|^2 + B|v|^2)}{\sqrt{q^2 + 4(A|u|^2 + B|v|^2)}}\right) = 0.$$

Saturable nonlinearity

The Lagrangian for these equations is

$$L = ir (\bar{u}u_t - u\bar{u}_t) + ir (\bar{v}v_t - v\bar{v}_t) - D_1 r |u_r|^2 - D_2 r |v_r|^2 + r \sqrt{q^2 + 4(A|u|^2 + B|v|^2)} - rq.$$

Next steps:

Insert trial functions into Lagrangian.

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Next steps:

- Insert trial functions into Lagrangian.
- Average Lagrangian $\int_{-\infty}^{\infty} L = 0.$
- \Box Variational equation \rightarrow Modulation Equations.
- Plus various conservation equations.

Nematicon Equations

Trial Functions

$$u = a_{1} \operatorname{sech} \left(\frac{\sqrt{(x-\xi_{1})^{2}+y^{2}}}{w_{1}} \right) e^{i\sigma_{1}+iV_{1}(x-\xi_{1})} + ig_{1}e^{i\sigma_{1}+iV_{1}(x-\xi_{1})}$$
$$v = a_{2} \operatorname{sech} \left(\frac{\sqrt{(x-\xi_{2})^{2}+y^{2}}}{w_{2}} \right) e^{i\sigma_{2}+iV_{2}(x-\xi_{2})} + ig_{2}e^{i\sigma_{2}+iV_{2}(x-\xi_{2})}.$$

Alternative trial functions

$$u = a_1 e^{-\frac{(x-\xi_1)^2+y^2}{w_1^2}} e^{i\sigma_1+iV_1(x-\xi_1)} + ig_1 e^{i\sigma_1+iV_1(x-\xi_1)}$$
$$v = a_2 e^{-\frac{(x-\xi_2)^2+y^2}{w_2^2}} e^{i\sigma_2+iV_2(x-\xi_2)} + ig_2 e^{i\sigma_2+iV_2(x-\xi_2)}.$$

Nematicon Equations

The first terms represent varying soliton-like pulses.

Second terms represent low wavenumber, low phase-speed radiation shed by the pulses. Should form circular shelves under the pulses.



where

 $\blacksquare a_k$: amplitude

- $\blacksquare w_k$: width
- ξ_k : displacement from x

 $\bullet \sigma_k$: phase

 V_k :

 $\blacksquare g_k$: shelf amplitude

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Radiation Loss

Shed radiation has small amplitude relative to the nematicons.

Governed by linearised NLS equations

$$iu_{z} + \frac{D_{1}}{2r} \frac{\partial}{\partial r} (ru_{r}) = 0$$
$$iv_{z} + \frac{D_{2}}{2r} \frac{\partial}{\partial r} (rv_{r}) = 0$$

Match shed radiation to existing radiation value at edge of shelf.

Mass lost to the shed radiation important.

- **Mass** loss and radiation shelf are both initially 0.
- Add mass loss term to equation for g_i .

















Two-colour nemations governed by CNLS-like equations.

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